

Optical Properties of Infrared Detector Materials

Stefan Zollner

Department of Physics, New Mexico State University, Las Cruces, NM

Email: zollner@nmsu.edu. WWW: <http://femto.nmsu.edu>.

Annual meeting of the
AVS Spectroscopic Ellipsometry Technical Group,
68th International AVS Symposium and Exhibition
Pittsburgh, PA, November 8th, 2022.



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Office of Scientific Research under
award number FA9550-20-1-0135.

Outline

1. Optical and X-ray Characterization of thick **GeSn alloys** (1-2% Sn) **on GaAs**
Ken Hass Outstanding Student Paper Award March 2022, Forum on Industrial & Applied Physics, APS
2. Coherent Longitudinal Acoustic **Phonon Oscillations** in Ge
using **Femtosecond** Pump-Probe Spectroscopic Ellipsometry (at ELI Beamlines, Prague, CR)
*phys. stat. solidi RRL **16**, 220058 (2022)*
3. Temperature Dependence of the **Infrared Dielectric Function of InSb** near the Band Gap
Two invited talks, 2023 DPG Spring and MRS Fall Meetings, submitted to JVST B

Maybe Next Year

4. *Structural Properties and Optical Constants of **CaF₂** from 30 meV to 9 eV (UG/MS project)*
5. **Photoluminescence** of Ge-Si-Sn Alloys under ambient and **high pressures** at low temperature
6. Spectroscopic Ellipsometry (30 meV to 6.5 eV) **at 10 K using Recirculating Helium Cooler**

Optical and X-Ray Characterization of $\text{Ge}_{1-y}\text{Sn}_y$ alloys on GaAs

Haley B. Woolf,¹ **Matt Kim**,² Carola Emminger,¹ Carlos A. Armenta,¹ Stefan Zollner¹

1. Department of Physics, New Mexico State University, Las Cruces, NM

2. QuantTera, Scottsdale, AZ



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Ge_{1-y}Sn_y on GaAs by CVD

- Provided by Matt Kim (commercial reactor)
- Sn content: 1.2% and 2.6% (XRD)
- Very thick: 160 and 115 nm (SE-interference)
- Fully pseudomorphic, no relaxation (XRD), broad asymmetric mosaic streak
- XRD, Ellipsometry, luminescence

W-Wavelength Streak

M-Mosaic Spread (Relaxation Line)

Black line drawn from origin (Relaxed)

Dashed line drawn through substrate peak

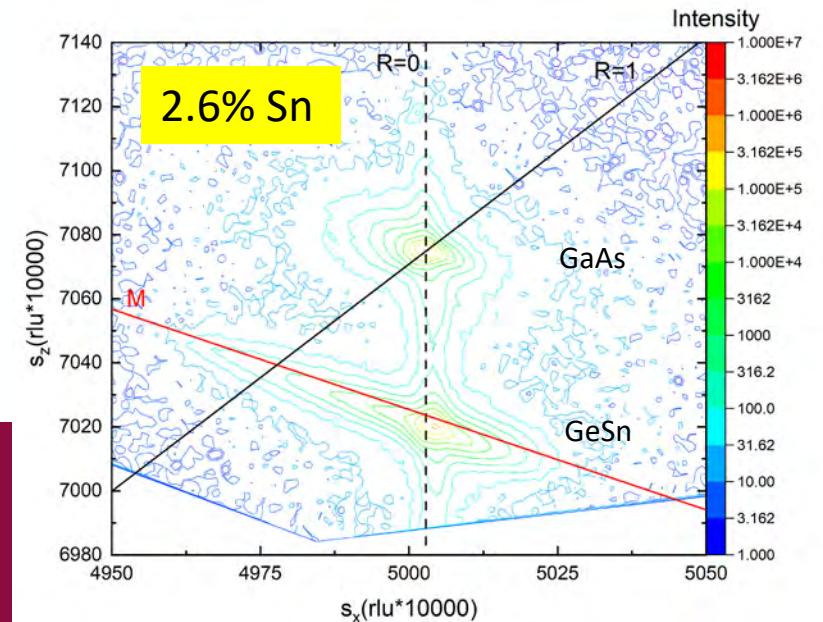
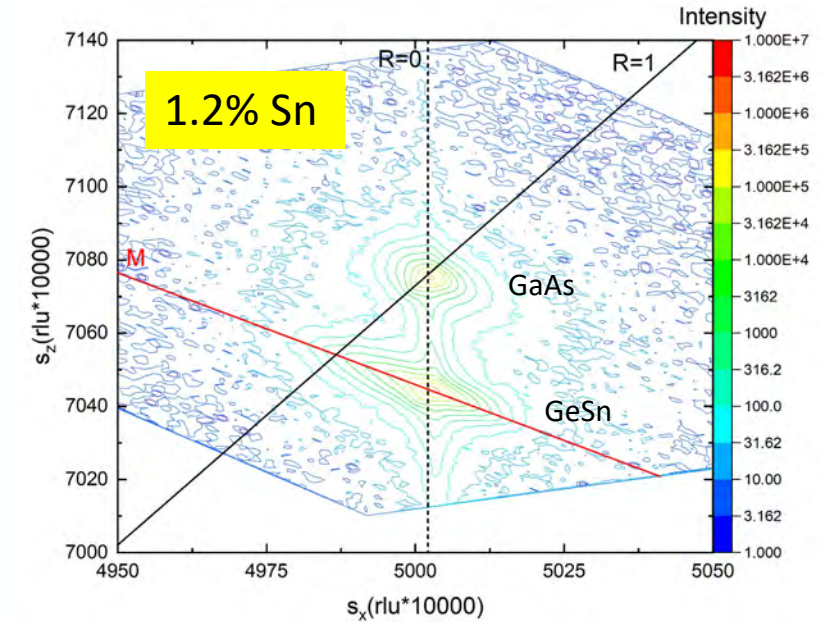
(Pseudomorphic)

$$s_x = \frac{q_x}{2\pi} = \frac{1}{\lambda} [\cos(\omega) - \cos(2\theta - \omega)]$$

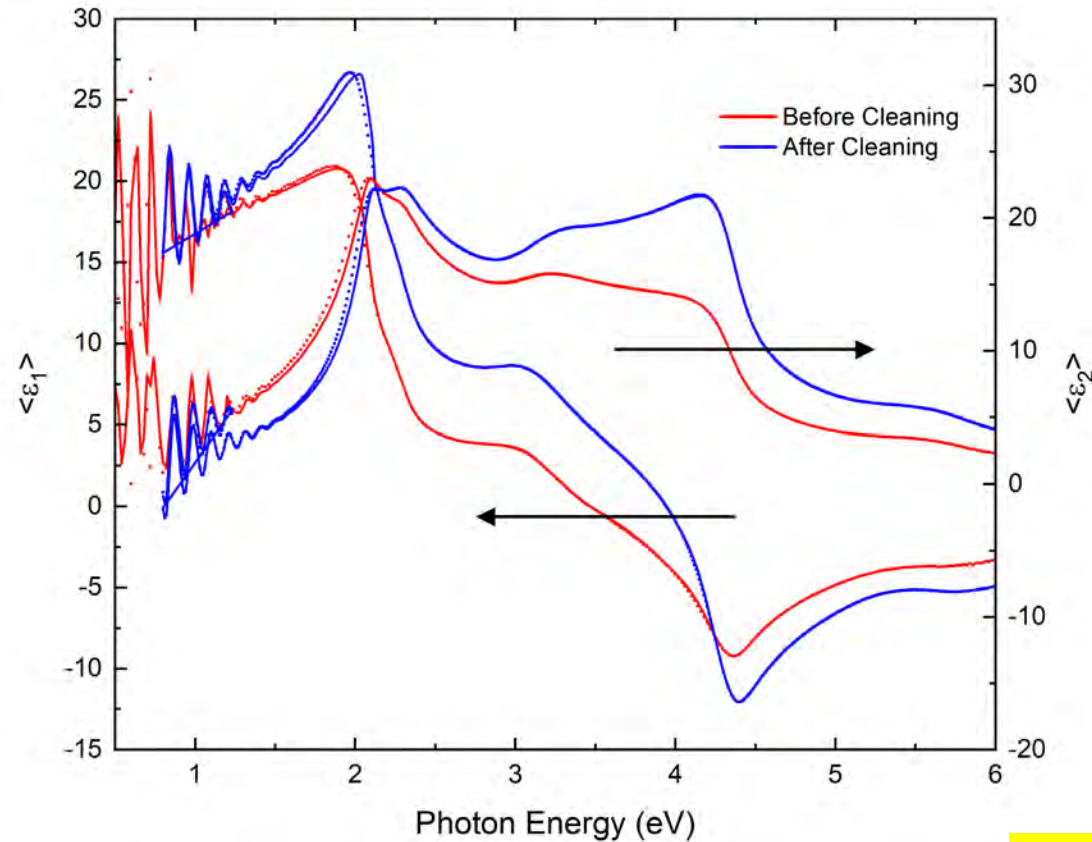
$$s_z = \frac{q_z}{2\pi} = \frac{1}{\lambda} [\sin(\omega) + \sin(2\theta - \omega)]$$

$$\lambda = 1.5406 \text{ \AA}$$

(224) g. i. RSM



Pseudodielectric Function Before/After Cleaning 2.6%

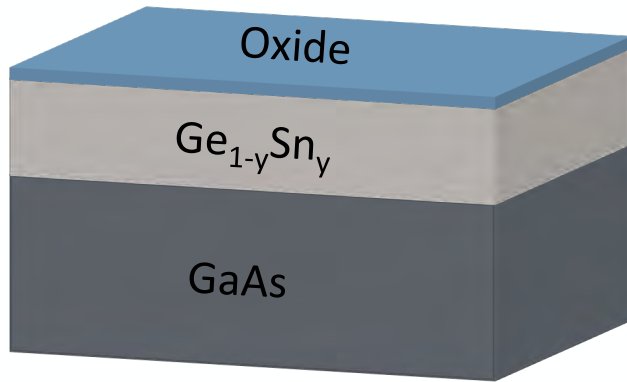


	Oxide Thickness
Before Cleaning (Oct 10 th):	68.27 Å
After Cleaning (Oct 11 th):	42.3 Å
After 2 nd Cleaning (Oct 14 th):	27.5 Å

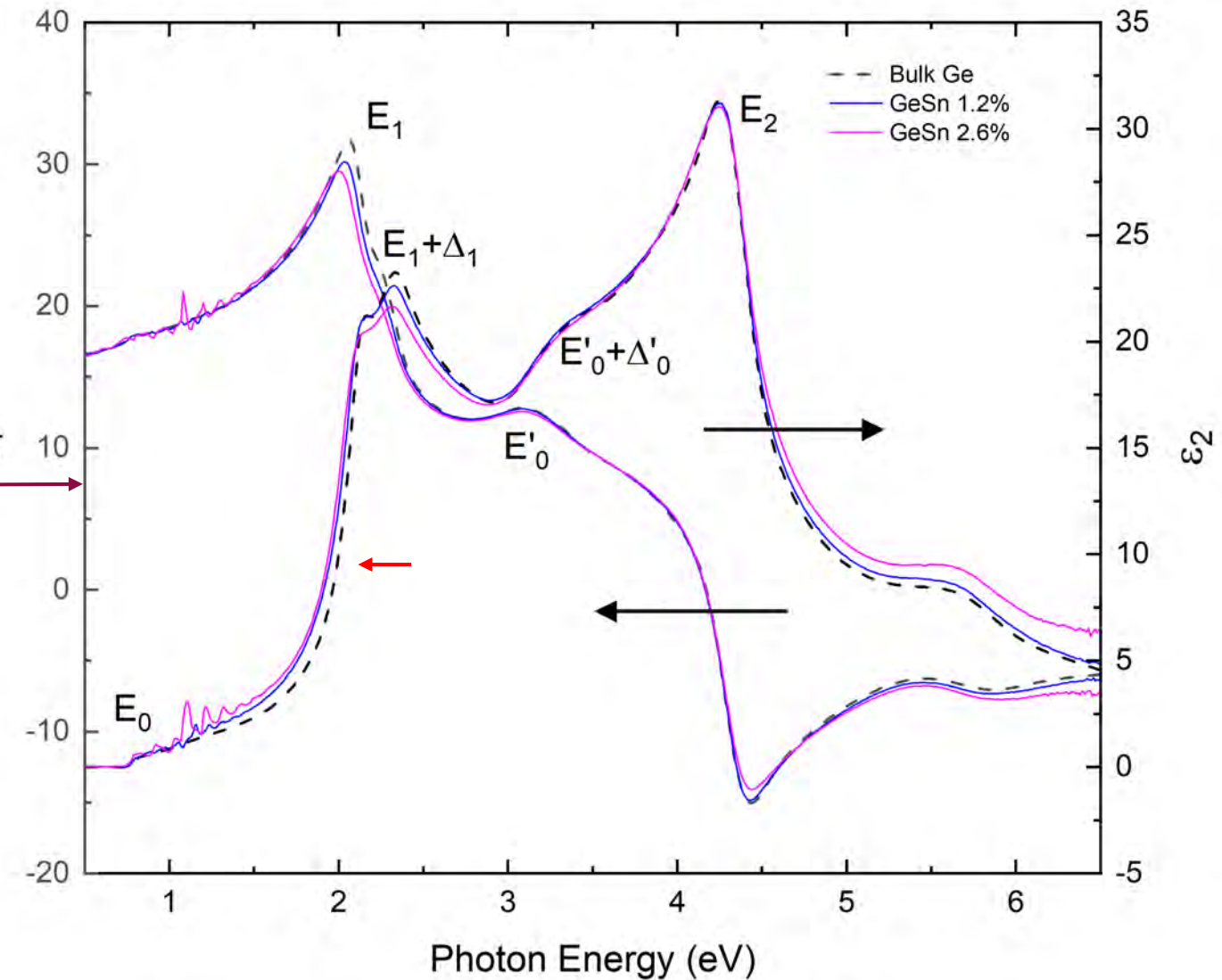
The $\text{Ge}_{1-y}\text{Sn}_y$ on GaAs sample was cleaned ultrasonically with **water** and then **isopropanol** for 10 mins each to remove organic layers and most of the native oxide.

The native oxide on Ge-Sn alloys is water soluble.

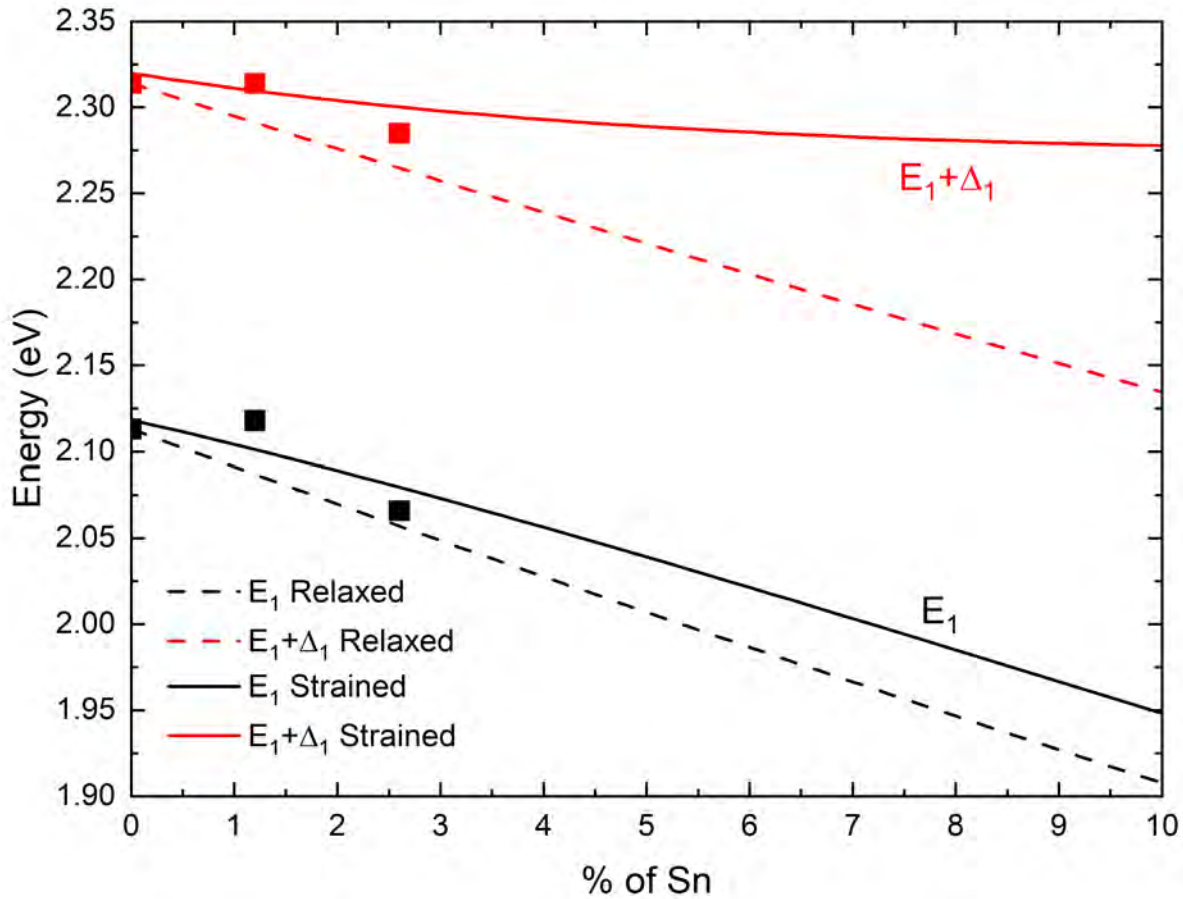
GeSn Optical Constants (after oxide correction)



- GeSn optical constants similar to Ge (low Sn content)
- Small redshift below E_1



Dependence of E_1 and $E_1 + \Delta_1$ on Tin Content



Relaxed:

$$E_1^{GeSn} = yE_1^{Sn} + (1 - y)E_1^{Ge} - b_{GeSn}y(1 - y)$$

$$(E_1 + \Delta_1)^{GeSn} = y(E_1 + \Delta_1)^{Sn} + (1 - y)(E_1 + \Delta_1)^{Ge} - b_{GeSn}y(1 - y)$$

$$b_{GeSn} = 1.350 \text{ eV}$$

Strained:

Continuum elasticity theory

$$E_1^{GeSn} = E_1^{Relaxed} + \Delta E_H - \frac{\Delta E_S^2}{\Delta_1}$$

small shear approximation
 $\Delta E_S \ll \Delta_1$

$$(E_1 + \Delta_1)^{GeSn} = (E_1 + \Delta_1)^{Relaxed} + \Delta E_H + \frac{\Delta E_S^2}{\Delta_1}$$

$$\Delta E_H = \sqrt{3}[yD_1^1^{Sn} + (1 - y)D_1^1^{Ge}]\epsilon_H$$

$$D_1^1^{Sn} = -5.4 \text{ eV}, \quad D_1^1^{Ge} = -8.7 \text{ eV}, \quad \epsilon_H = \frac{\epsilon_{\perp} + 2\epsilon_{\parallel}}{3}$$

$$\Delta E_S = \sqrt{6}[yD_3^3^{Sn} + (1 - y)D_3^3^{Ge}]\epsilon_S$$

$$D_3^3^{Sn} = -3.8 \text{ eV}, \quad D_3^3^{Ge} = -5.6 \text{ eV}, \quad \epsilon_S = \frac{\epsilon_{\perp} - \epsilon_{\parallel}}{3}$$

More scatter than usual, but consistent with pseudomorphic strain.

Photoluminescence Data at 300 K (900 mW, 808 nm)

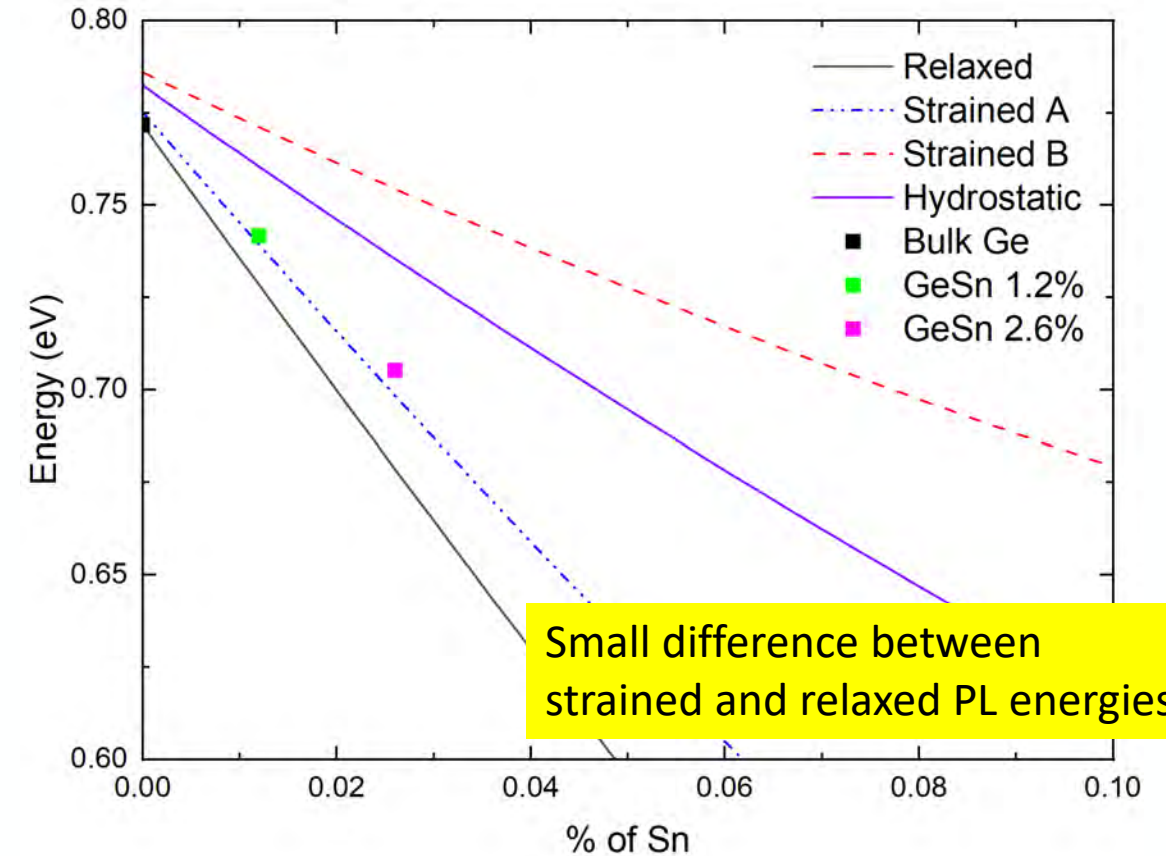
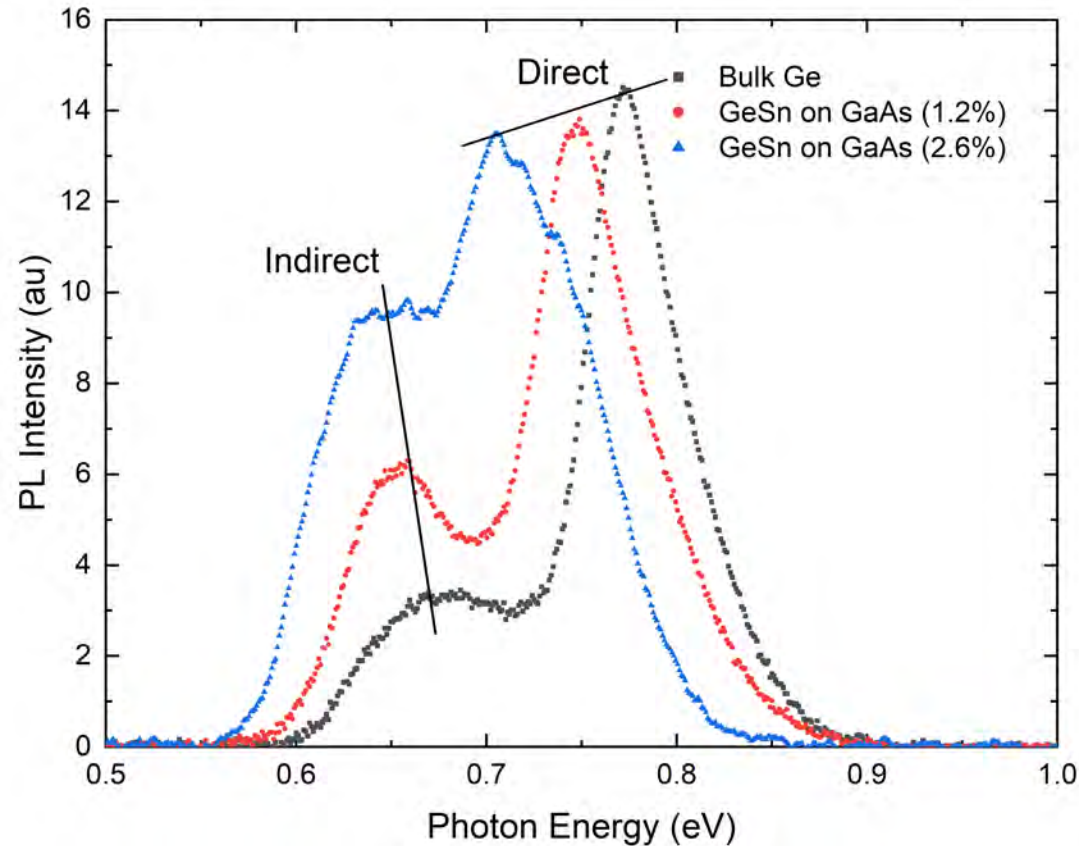
Relaxed:

$$E_0^{GeSn} = yE_0^{Sn} + (1 - y)E_0^{Ge} - b_{GeSn}y(1 - y)$$

$$b_{GeSn} = 2.46 \text{ eV}$$

Small Shear Strain:

$$E_0^{GeSn} = E_0^{Relaxed} + 3a\epsilon_H \pm \frac{3}{2}b\epsilon_S$$

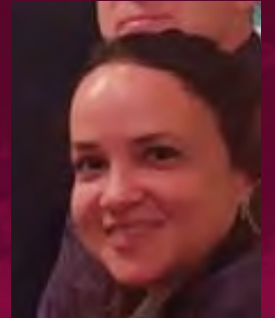


Haley can analyze Raman data under pressure with the same formalism.

Coherent Acoustic Phonon Oscillations in Ge Using Pump-Probe Time-Resolved Spectroscopic Ellipsometry

Carlos A Armenta,¹ Martin Zahradnik,² Carola Emminger,^{3,4} Shirly Espinoza,²
Mateusz Rebarz,² Jakob Andreasson,² Stefan Zollner¹

1. Department of Physics, New Mexico State University, Las Cruces, NM
2. ELI Beamlines, Dolní Brežany, Czech Republic
3. Department Physics, Leipzig University, Leipzig, Germany
4. Department of Physics, Humboldt University of Berlin

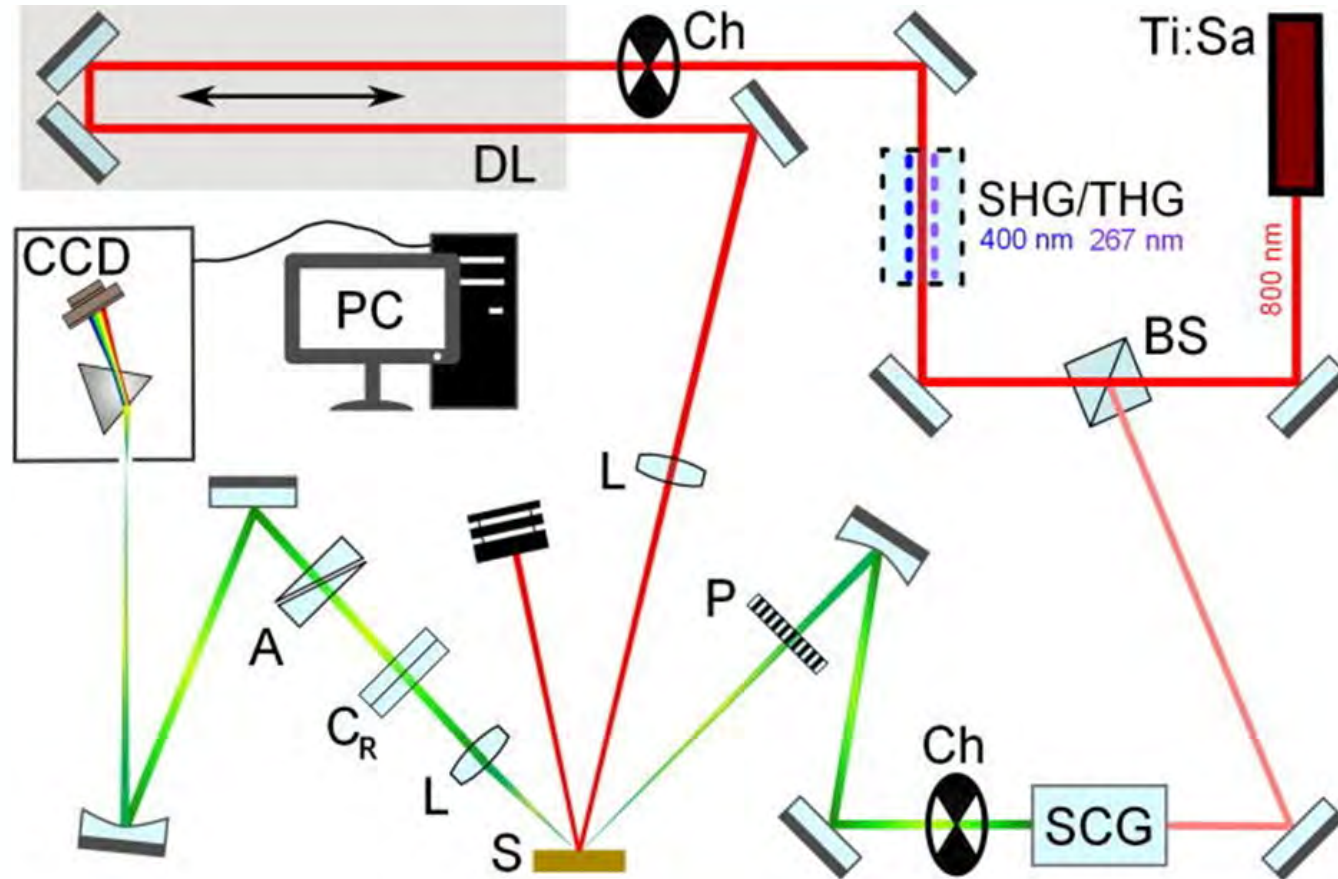


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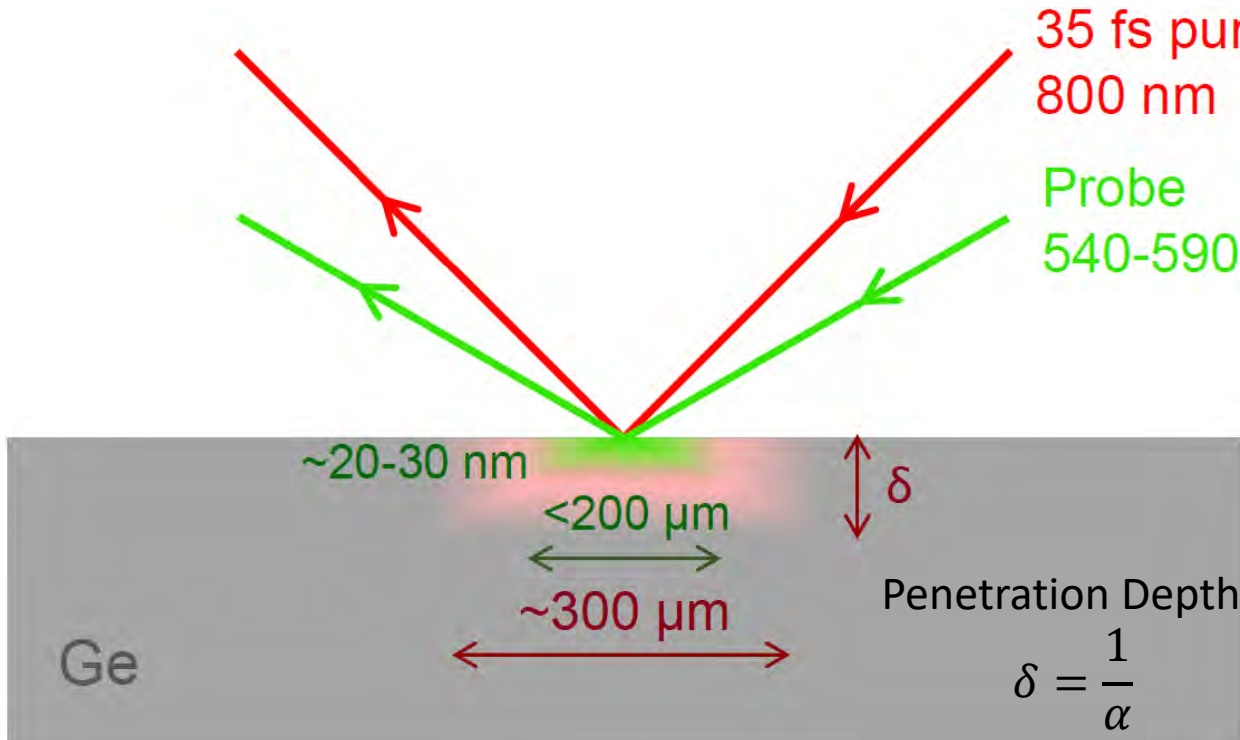
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Experimental Setup: Femtosecond Pump-probe Ellipsometry

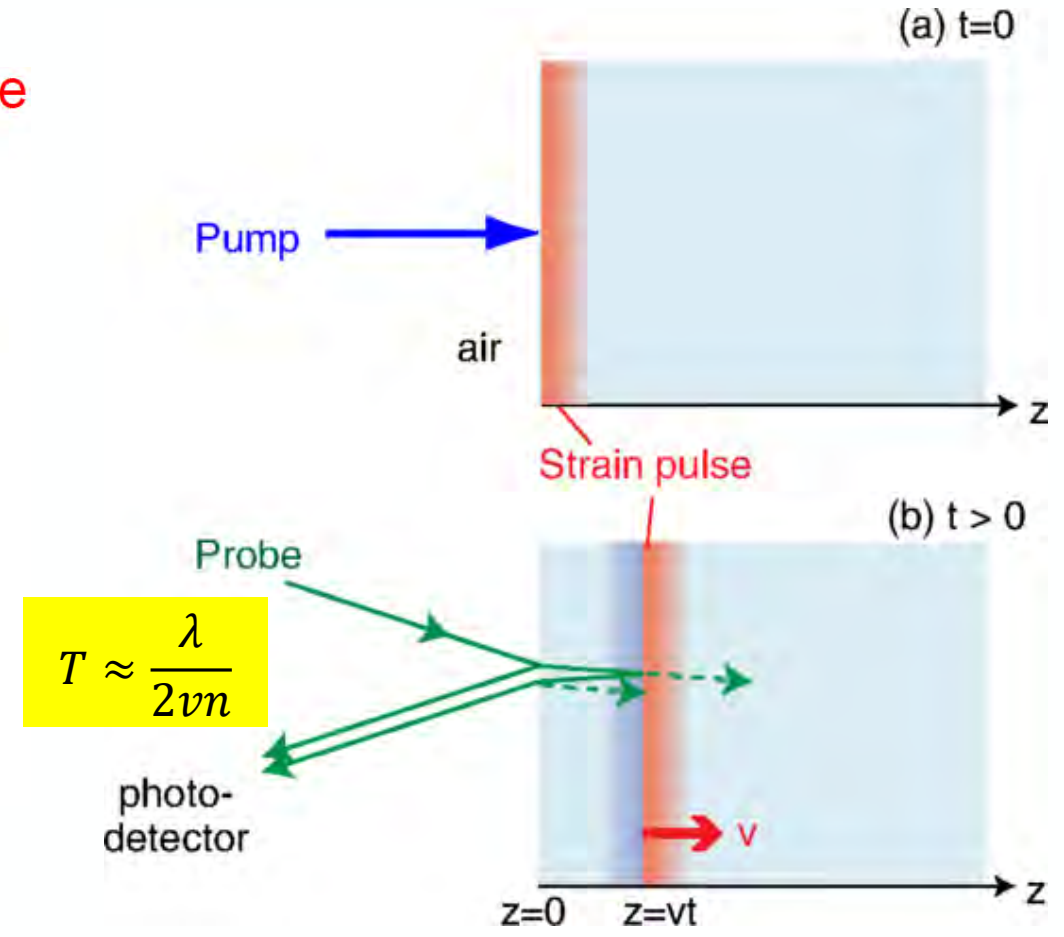


- Ch: Chopper (500 Hz, 250 Hz)
- A: Analyzer
- P: Polarizer
- C_R : Rotating Compensator
- L: Lens
- S: Sample
- DL: Delay Line (~ 6.67 ns pump-probe delay and 3 fs resolution)
- BS: Beam Splitter
- SHG/THG: Second/Third Harmonic Generation
- SCG: Super-continuum Generation
- CCD: Charge-coupled device detector

Femtosecond pump-probe ellipsometry



Carrier concentration: $1-3 \times 10^{21} \text{ cm}^{-3}$
Band filling and exciton screening

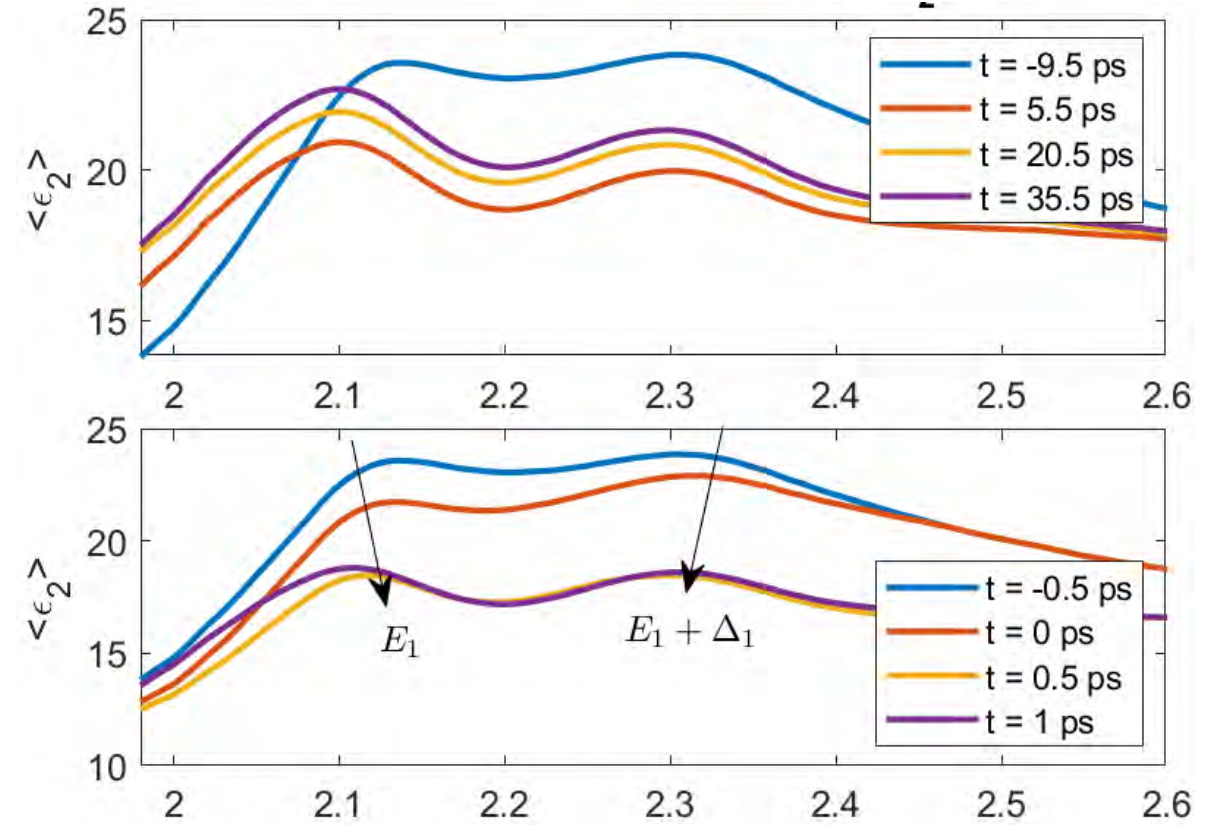
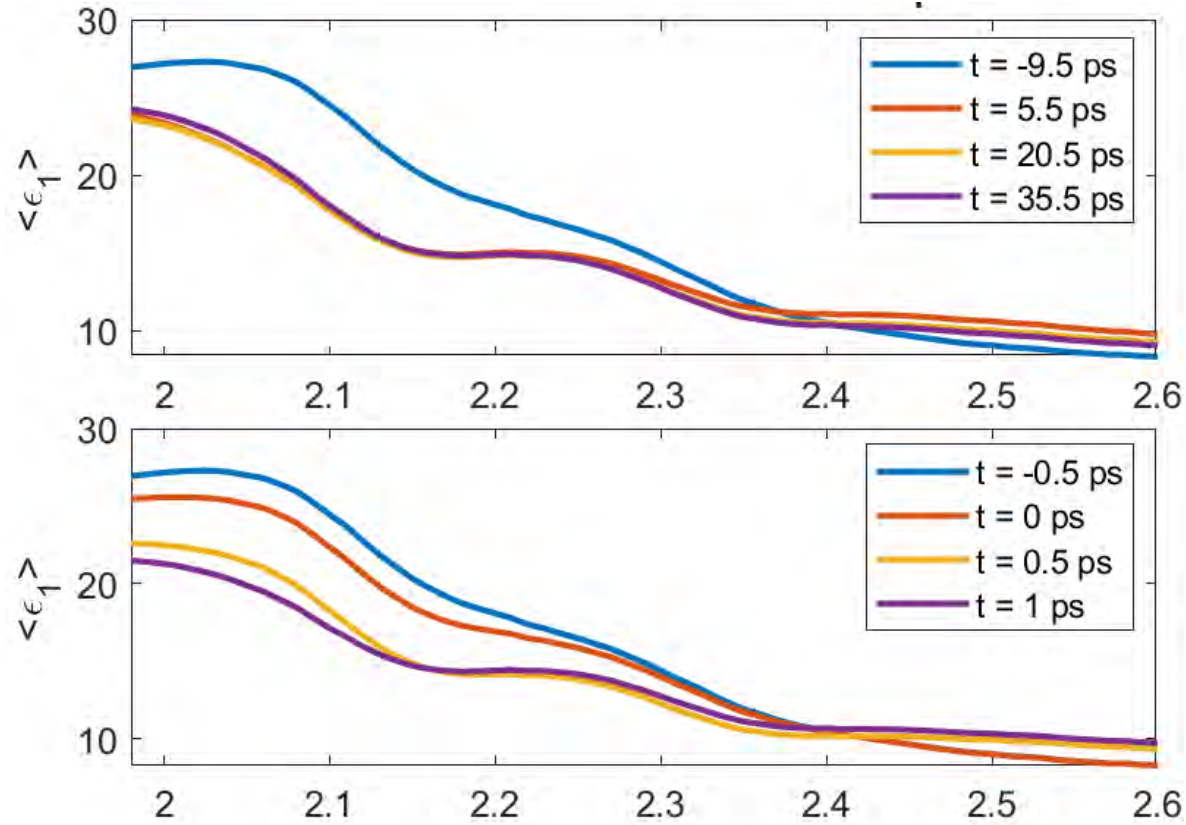


Coherent longitudinal acoustic phonon oscillations



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Pseudo-dielectric constant as function of delay time



E_1 energy and broadening change as a function of time.
Absorption is reduced and recovers.
Band filling, exciton screening, band gap renormalization.
Modeling is in progress: C. Xu, JAP **125**, 085704 (2019)

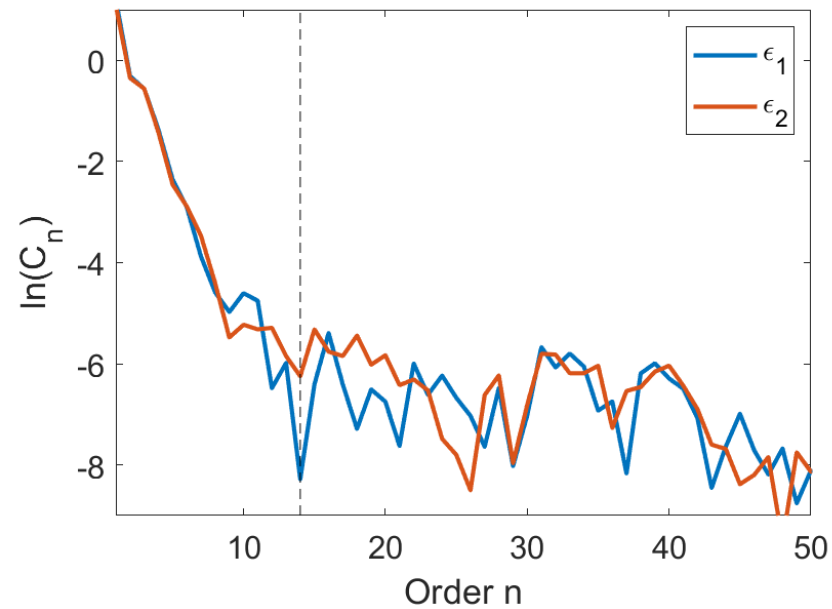
High Intensity Ge (100)
 $3 \times 10^{21} \text{ cm}^{-3}$

Digital Filtering with Fourier transforms, Gauss filters

Reduce noise

Calculate second derivative of ellipsometry spectra

Fourier Coefficients C_n



Direct Space Convolution

$$\bar{f}(E) = \int_{-\infty}^{\infty} dE' f(E') b_M(E - E')$$

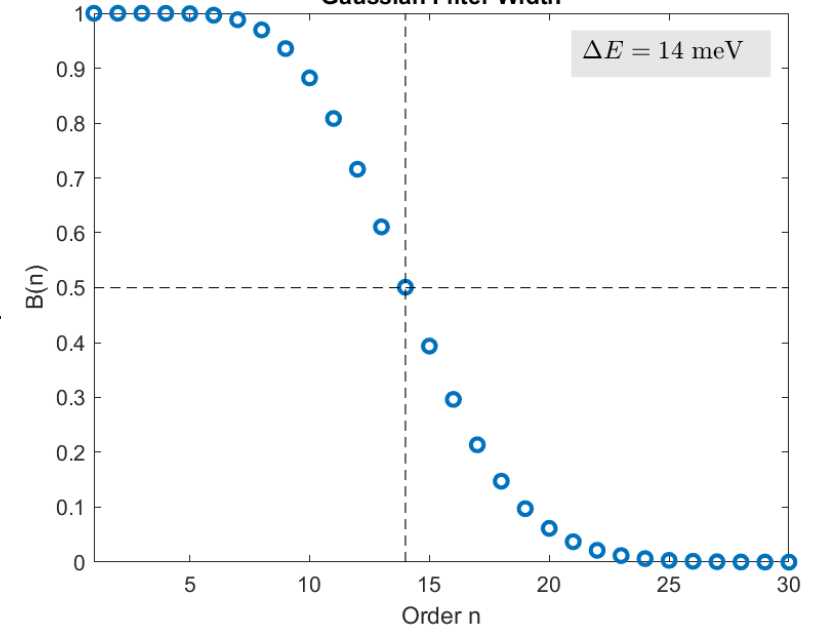
Extended Gauss (EG) Filter:

$$b_M(x) = \sum_{M=0}^{\infty} (-1)^M \frac{a^M}{M!} \frac{d^M}{da^M} \frac{1}{2\sqrt{\pi a}} e^{-\frac{x^2}{4a}}$$

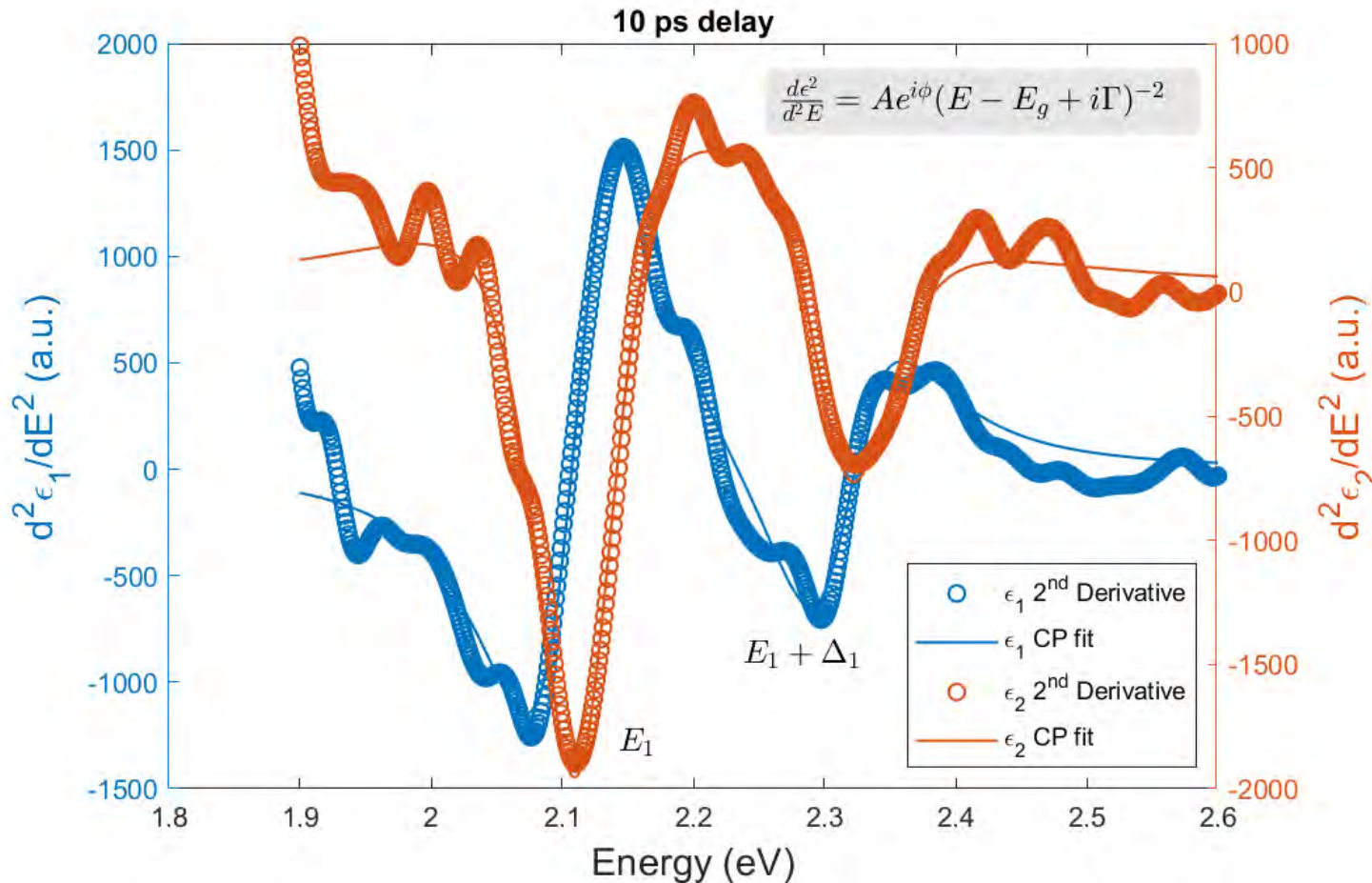
$$a = \frac{1}{\Delta k^2} = \Delta E^2$$

Filter width ΔE is chosen according to white noise onset.

Gaussian Filter Width



Critical Point Parameters as a Function of Delay Time



E_1 and $E_1 + \Delta_1$ critical points:

Second derivative calculated using linear filters method (LKKA).

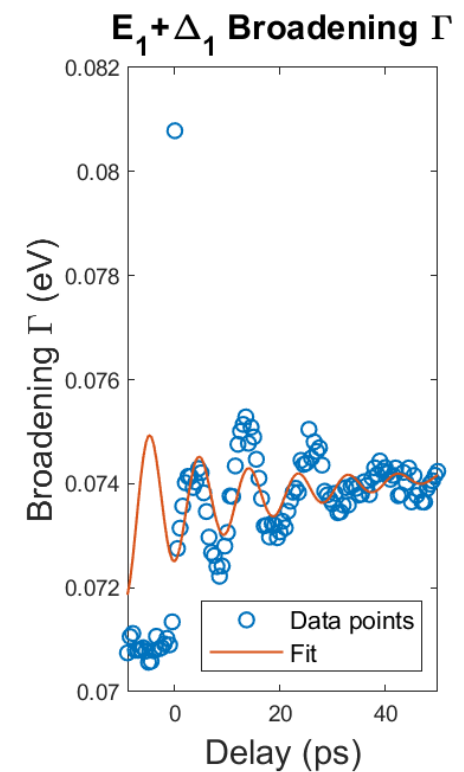
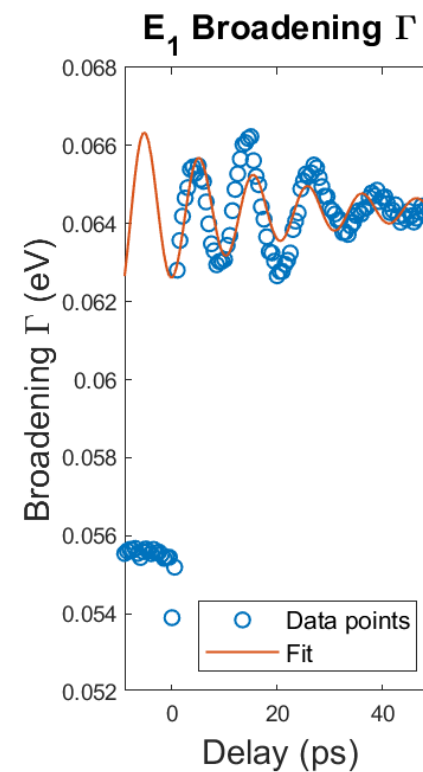
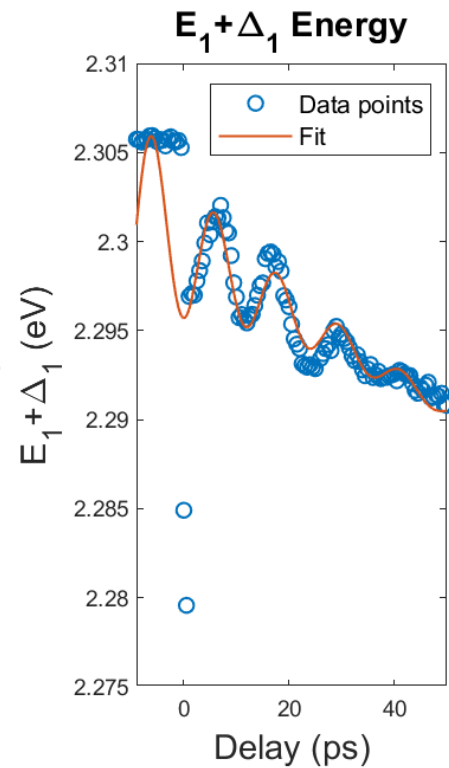
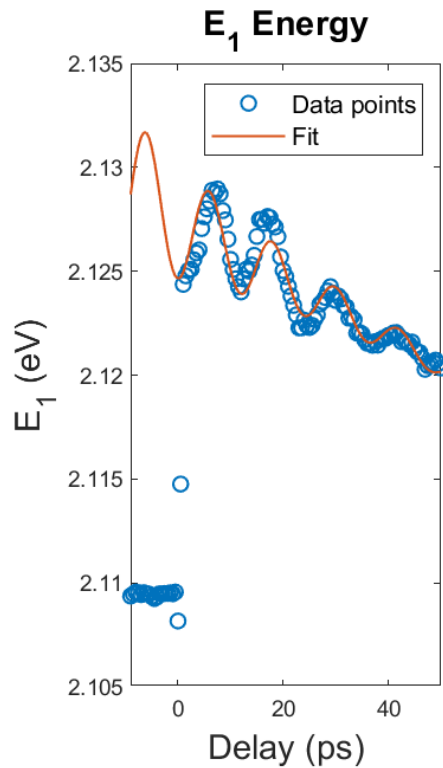
Model dielectric function

$$\epsilon_{2D}(E) = B - Ae^{i\phi} \ln(E - E_g + i\Gamma),$$

where:

- **Amplitude A**
- **Phase angle ϕ**
- **Energy E_g**
- **Broadening Γ**

E_1 and $E_1 + \Delta_1$ CP parameters: Acoustic phonon oscillations



Oscillations fitting function:

$$f(t) = A \cos(\omega t) e^{-ct} - bt + c$$

$$T \approx \frac{\lambda}{2vn}$$

Dependence on surface orientation is under investigation.

High Intensity Ge (100)
11/26/2021*



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Compare:

Carola Emminger, phys. stat. solidi RRL 16, 220058 (2022)

Period of oscillation of parameters

$$T \approx \frac{\lambda}{2vn}$$

Oscillation Period						
	(111) HI	(111) MI	(110) HI	(110) MI	(100) HI*	(100) MI
Theory Period E_1 (ps)	9.39		9.65		10.62	
Theory Period $E_1+\Delta_1$ (ps)	9.65		9.92		10.91	
Energy E_1 (ps)	12.2 (0.6)	11.5 (1.2)	9.2 (0.9)	8.6 (1.3)	12.0 (0.4)	10.7 (0.8)
Energy $E_1+\Delta_1$ (ps)	14.6 (3.3)	11.0 (1.1)	9.5 (1.5)	8.4 (0.7)	11.9 (0.4)	11.0 (0.5)

Good agreement with theory for (100) surface orientation.

Dependence on surface orientation is under investigation.

Error bars too large to see dependence of period on surface orientation.

More measurements in February 2023.

Temperature Dependence of the Infrared Dielectric Function of InSb near the Direct Band Gap



FA9550-20-1-0135
W911NF-2210130

MELISSA RIVERO ARIAS, CAROLA EMMINGER, CARLOS ARMENTA, NUWANJULA SAMARASINGHA,
CESY M. ZAMARRIPA, JADEN R. LOVE, AND STEFAN ZOLLNER

Department of Physics
New Mexico State University, Las Cruces, NM, USA

Submitted to J. Vac. Sci. Technol. B (under review)

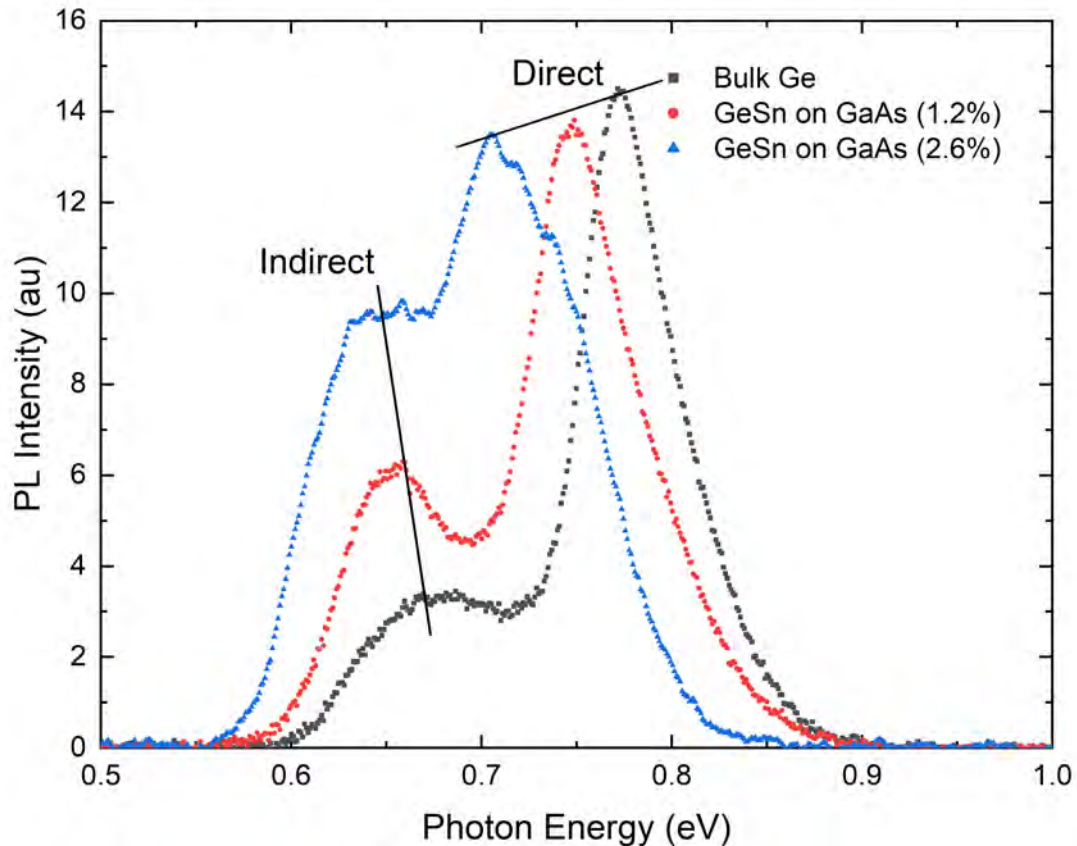
Email: zollner@nmsu.edu. WWW: <http://femto.nmsu.edu>.



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Motivation: Modeling of Detector Spectral Response

Photoluminescence of Ge-Sn Alloys:
300 K, 900 mW, 808 nm



Questions:

- If we see a photoluminescence spectrum, can you calculate the detector response?
- Can we calculate the absorption spectrum?
- Would that be useful?

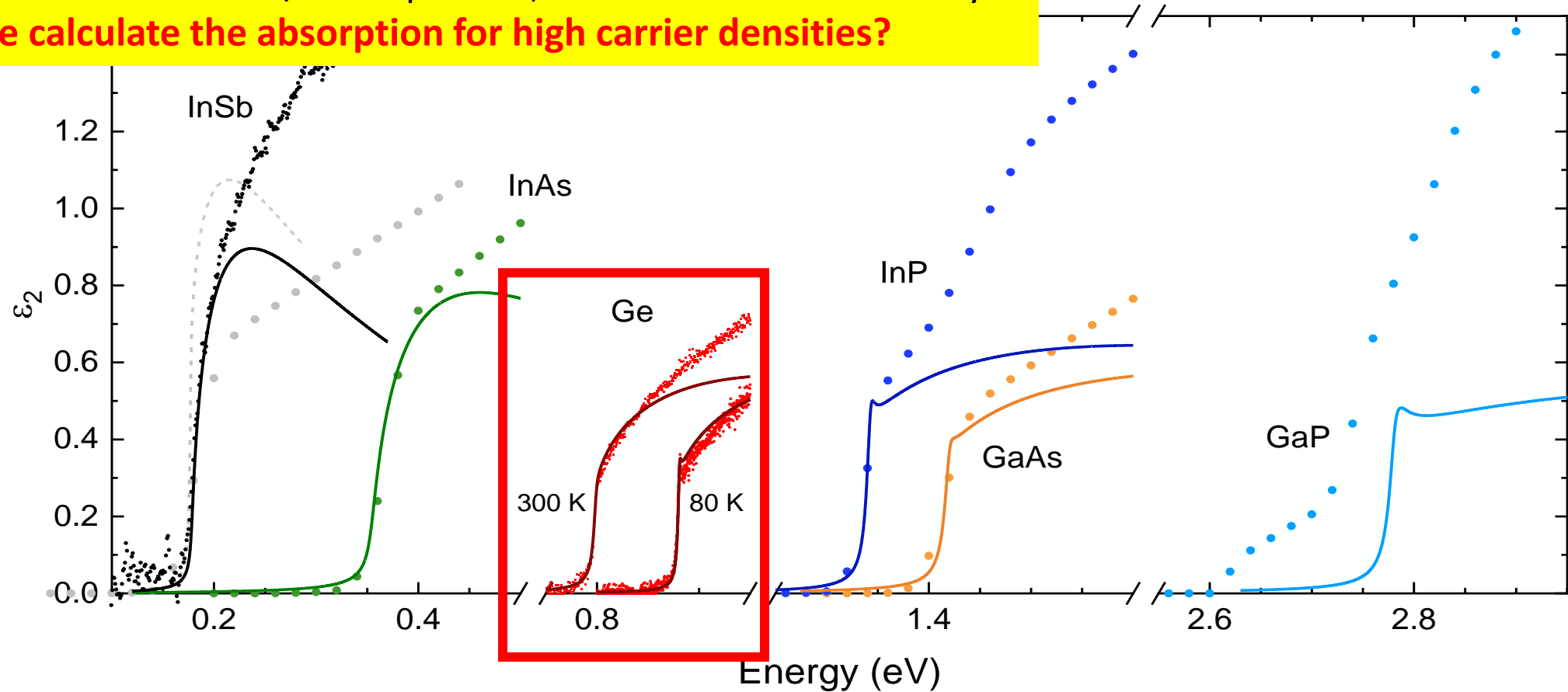
Calculation of Absorption Spectrum from k.p Theory

Can we calculate the absorption spectrum?

Yes, we can for Ge in the low carrier density limit.

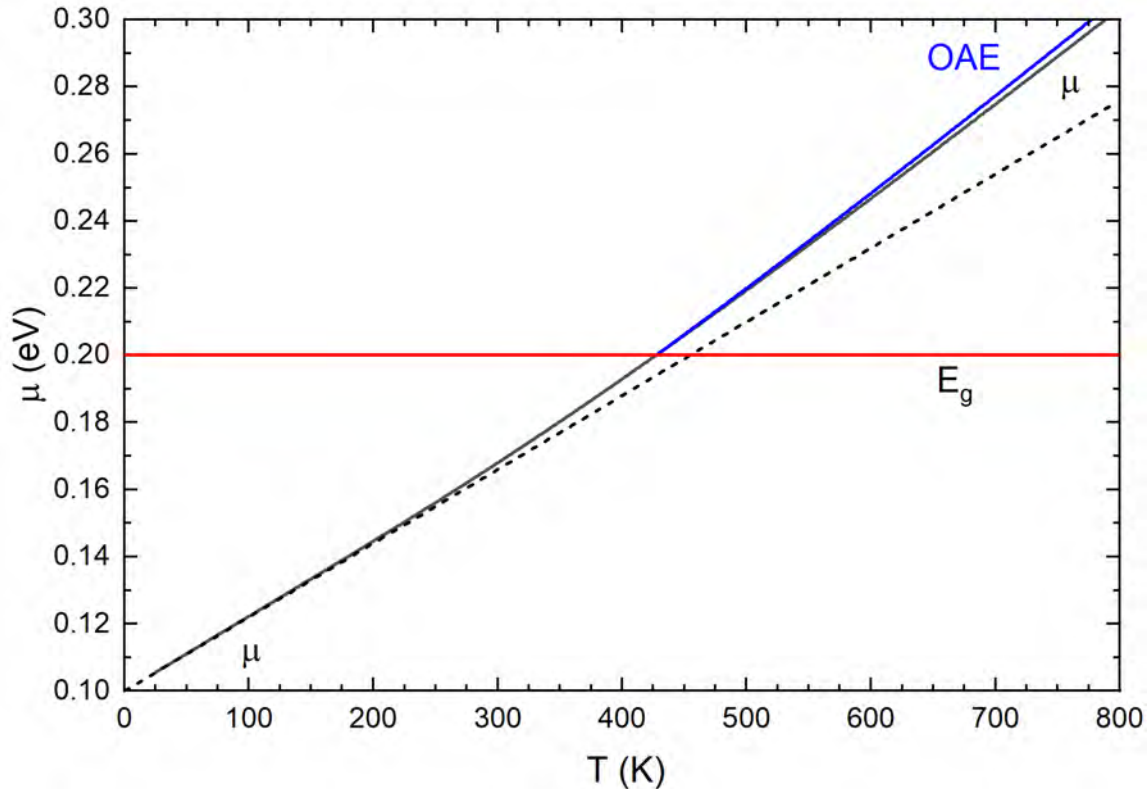
It does not work for other III/V compounds, have not tried Ge-Sn alloys.

How can we calculate the absorption for high carrier densities?

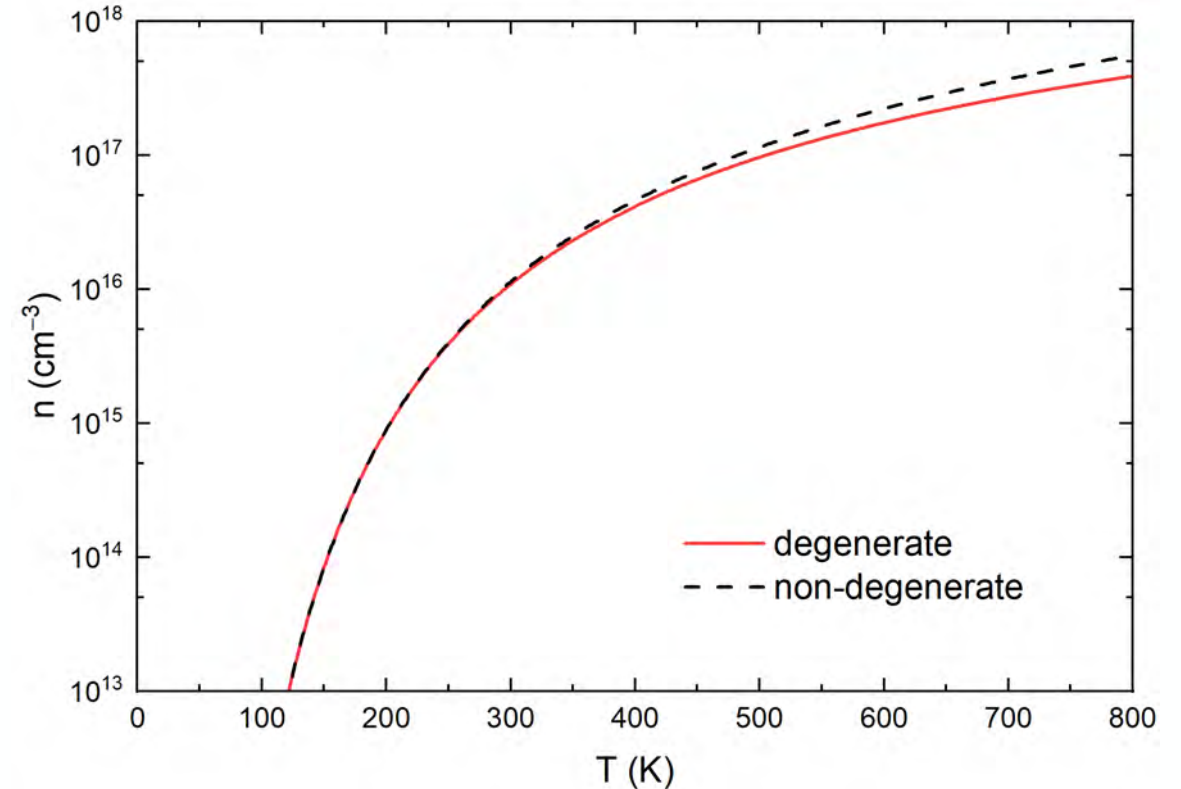


Temperature Changes Intrinsic Carrier Concentration in InSb

Above 400 K, the Fermi level is above the conduction band minimum.

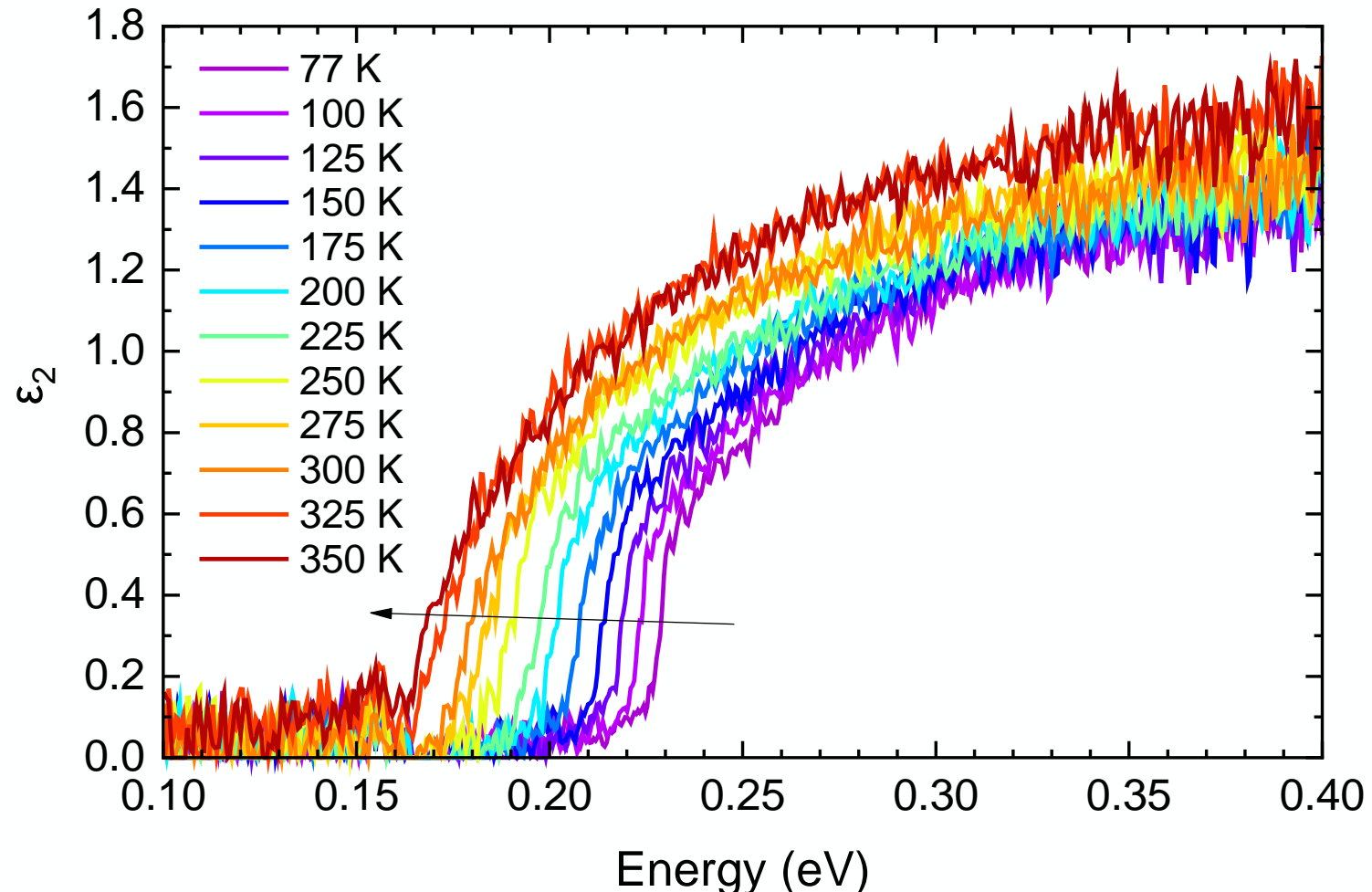


Carrier concentration reaches 10^{18} cm^{-3} near the melting point of InSb (800 K)



InSb: How do many-body effects influence the absorption by MWIR detectors?

InSb (100): Initial attempt promising, but that is not the whole story

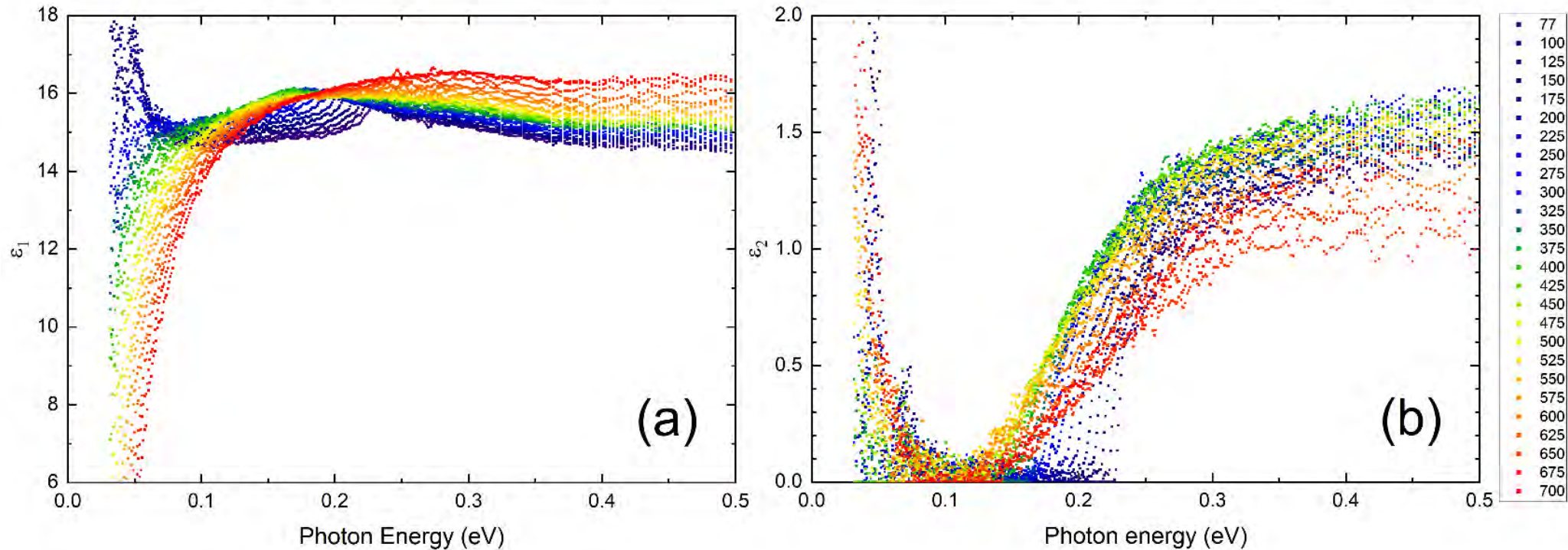


Multiple experimental issues:

- InSb sample cracks, melts, reacts with the Cu sample holder.
- Adhesive (carbon nanoparticles, silver paint) expands, evaporates, redeposits on the windows.
- Beam larger than sample: Depolarization from sample holder reflections.
- Black-body radiation, heat shield.
- Cryostat leaks, thermocouple breaks.

Initial result: **Clear redshift with increasing temperature (up to 450 K).** Strange things happen above 450 K.

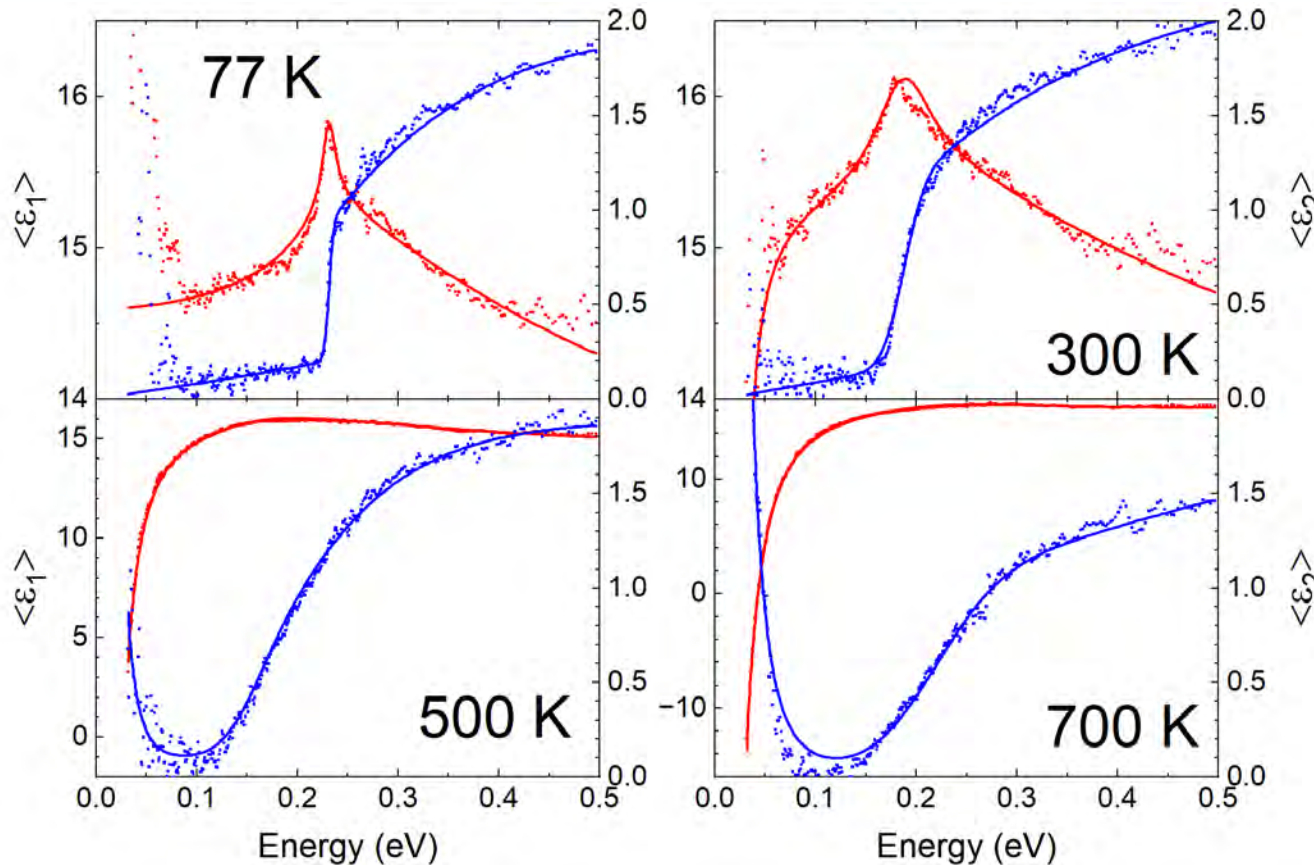
Dielectric Functions of InSb from 80 to 800 K



- Band gap changes with temperature (but only below 500 K).
- Reduction of absorption coefficient at high temperatures.
- Drude response at high temperatures (thermally excited carriers).
- *Depolarization artifacts at long wavelengths (below 300 K).*

Band gap analysis: Parametric semiconductor model

- How does the band gap of InSb change with temperature?



Parametric-Semiconductor Model:

Parameterized Semiconductor Layer

Layer Name: PSEMI

Comment: Parameterized Semiconductor Layer

Thickness: 1 mm

Position (eV): Pole #1: 8, Pole #2: 0.02

Magnitude: 3.2463, 1e-005

Optical Constants >>

Opt Conat Fit: n, k

Buttons: Ok, Delete Layer, Replace Layer

Joint DOS Parameters: Change

Set	Energy	Amp	Connect	Br	Disc	Mid Pos	Mid Amp	2nd order	Mid Pos	Mid Amp	2nd order
R#:	0.2262	F	0.3141	F	0	2	4.748	F	0.9990	F	0.5000
R#:	0.3900	F	0.0783	F	1	2	45.000	F	0.0000	F	0.5000
R#:	1.8079	F	15.7720	F	0	4	56.682	F	0.2768	F	0.4519
R#:	2.3086	F	8.3773	F	0	4	61.667	F	0.1430	F	0.3000
R#:	3.5529	F	12.2446	F	3	8	177.396	F	-0.3500	F	0.4000
R#:	3.8727	F	47.1769	F	3	8	244.267	F	-0.8016	F	0.8000
R#:	5.2758	F	1.8163	F	3	8	250.000	F	-0.9500	F	0.8000
R#:	5.8715	F	1.0438	F	3	8	300.000	F	-0.9500	F	0.8000
R#:	7.0000	F	2.9256	F	7	8	700.000	F	0.0000	F	0.5000
R#:	4.5000	F	10.0000	F	8	10	50.000	F	0.0000	F	0.5000
R#:	5.0000	F	10.0000	F	9	11	50.000	F	0.0000	F	0.5000
R#:	5.5000	F	10.0000	F	10	12	50.000	F	0.0000	F	0.5000

Also vary “shape parameters”.

Asymmetric peak shape poorly described.

Try Tanguy oscillator for excitonic line shape.

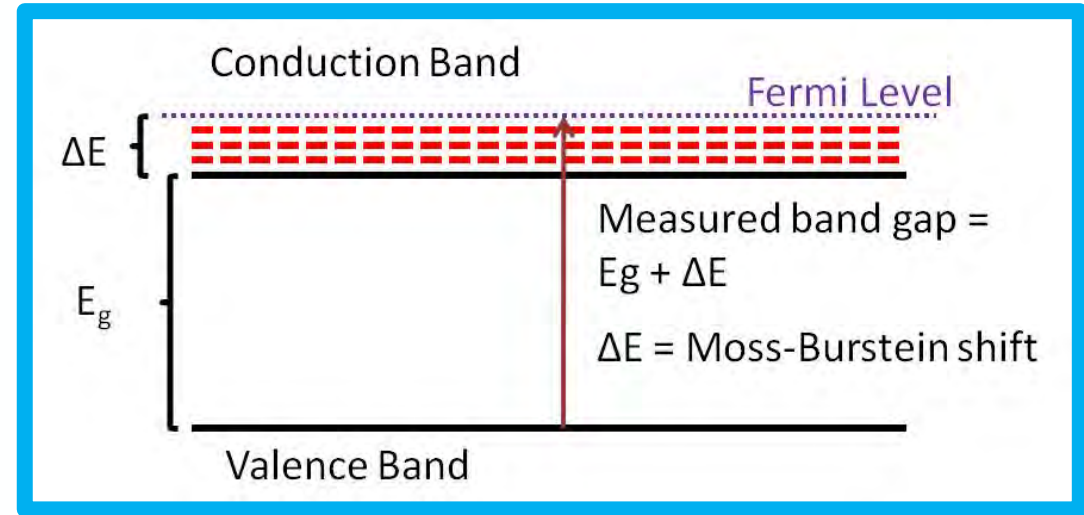
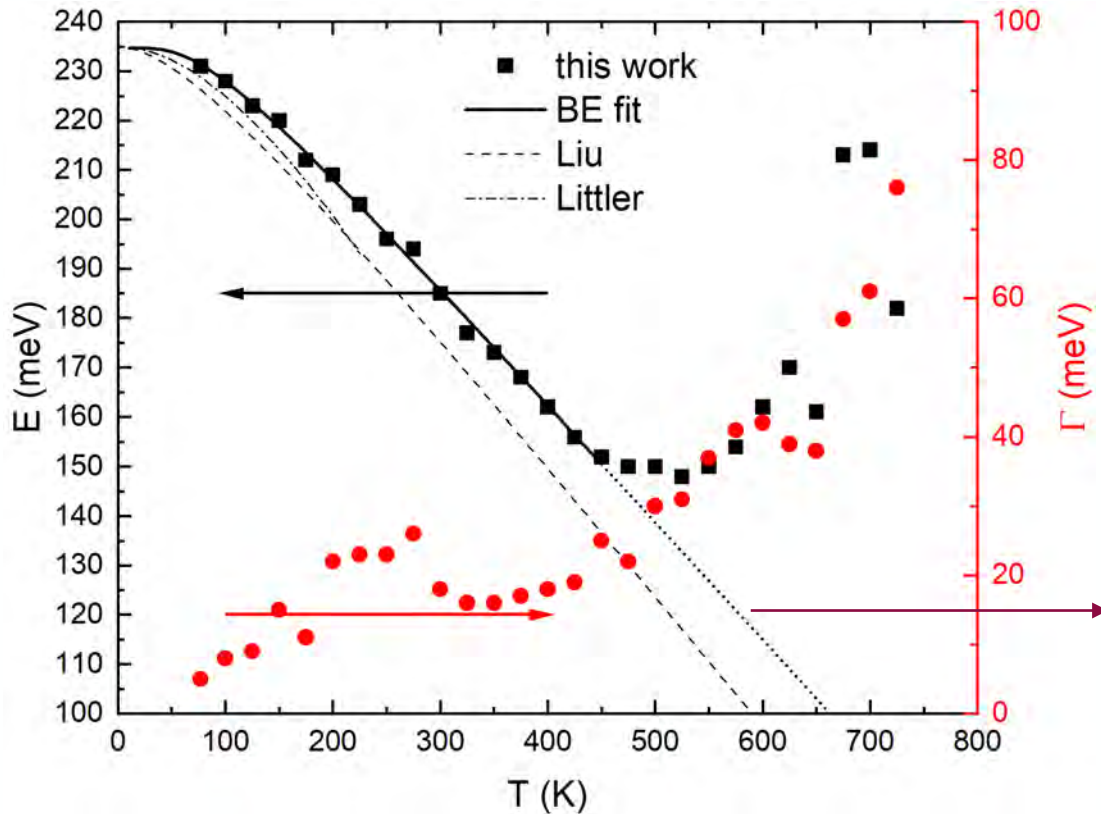
Fit	Final
MSE	0.2958
En0.0	0.22615 ± 0.000889
Br0.0	4.7478 ± 1.32
Am0.0	0.31415 ± 124
Disc0.0	0.999 ± 788
RPos0.0	0.84009 ± 0.0264
RAmp0.0	1.8912 ± 0.191
PoleMag.0	3.2469 ± 6.56
PoleMag2.0	1e-005 ± 0.000568



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C. M. Herzinger, B. Johs, et al., J. Appl. Phys. **83**, 3323 (1998).

Direct Band Gap of InSb versus Temperature



Bose-Einstein

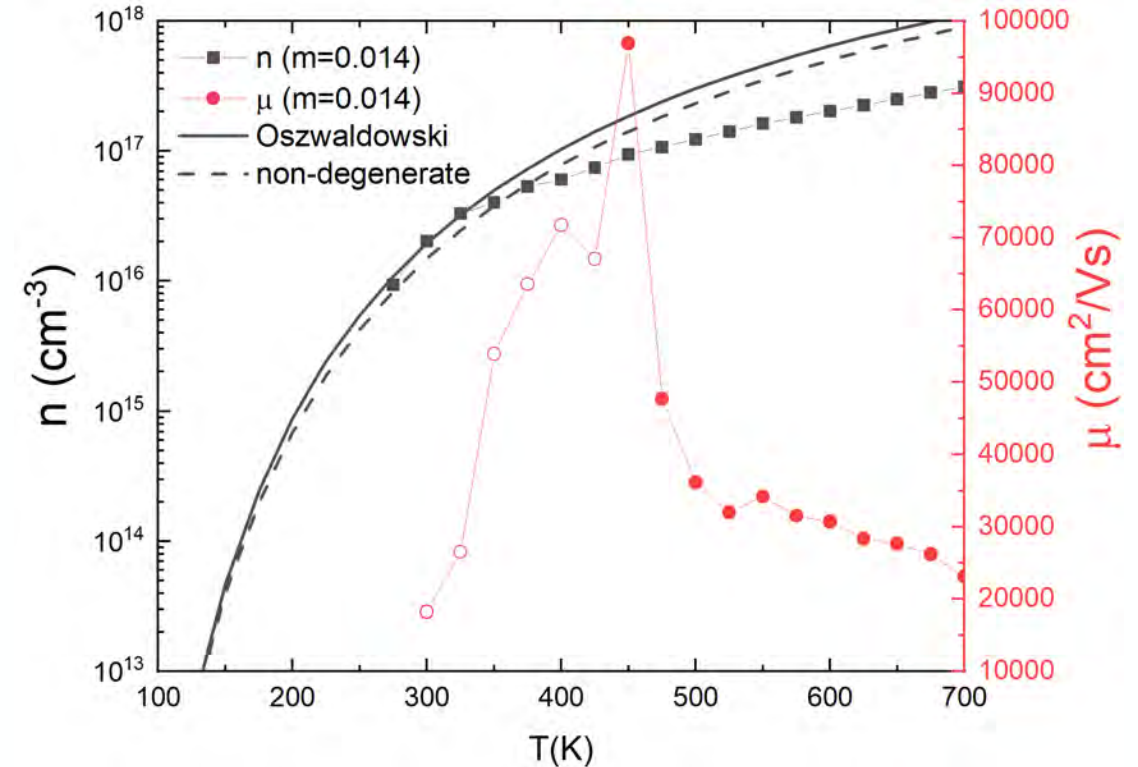
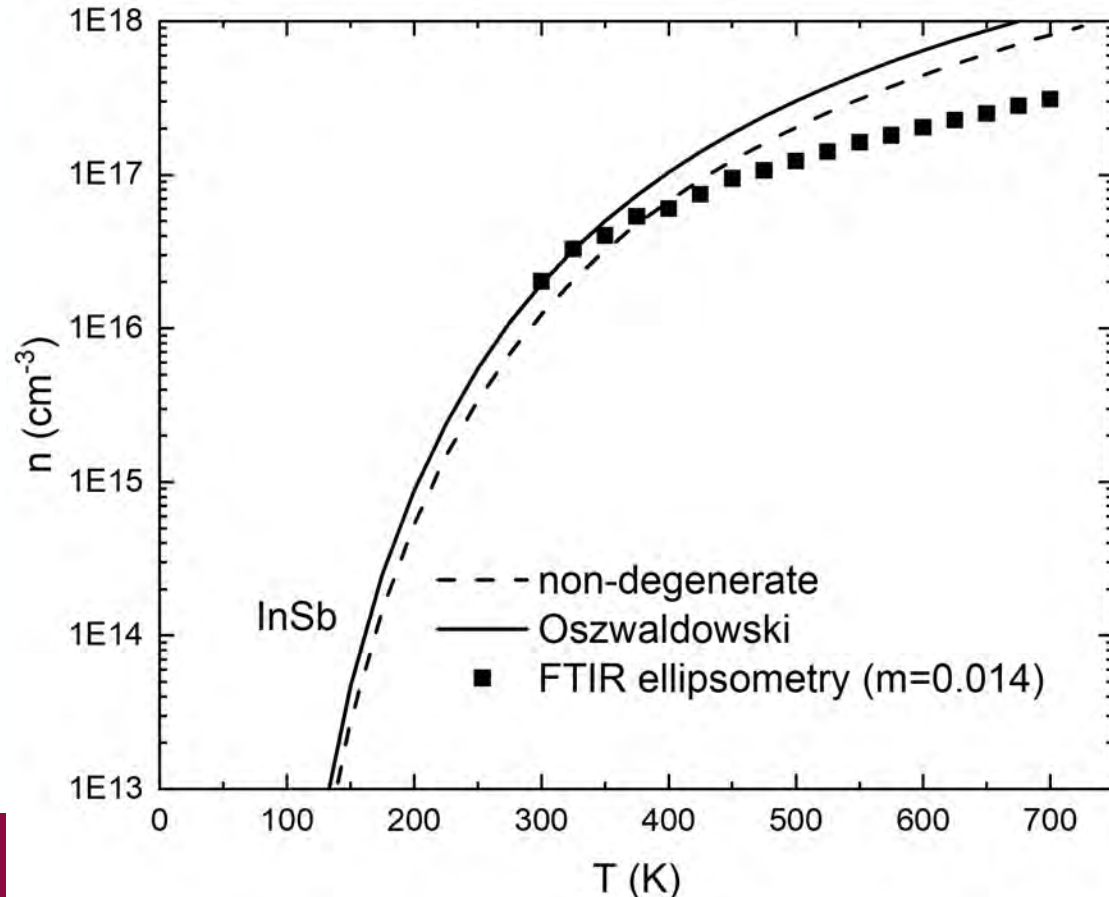
$$E_g(T) = a - b \left[1 - \frac{2}{e^{\bar{T}} - 1} \right]$$

- Band gap changes with temperature (but only below 500 K)
- Described by Bose-Einstein model below 500 K: Logothetidis, PRB **31**, 947 (1985).
- Increase above 500 K: thermal Burstein-Moss shift.

Free-carrier absorption

To model the Drude response, we make some really bad assumptions:

- Parabolic bands (ignore non-parabolicity)
- Effective mass constant $m_e=0.014$ (independent of temperature)



Oszwaldowski/Zimpel, J. Phys. Chem. Solids **49**, 1179 (1988).
D. L. Rode, Phys. Rev. B **3**, 3287 (1971).

Required model improvements

$$\varepsilon_2(E) = \frac{2\pi A\sqrt{R}}{E^2} \left\{ \sum_{n=1}^{\sqrt{g}} \frac{2R}{n} \left(\frac{1}{n^2} - \frac{n^2}{g^2} \right) \delta \left[E - E_0 + \frac{R}{n^2} \left(1 - \frac{n^2}{g} \right)^2 \right] + \frac{\sinh(\pi g k) H(E - E_0)}{\cosh(\pi g k) - \cosh \left(\pi g \sqrt{k^2 - \frac{4}{g}} \right)} \right\} [f_h(E) - f_e(E)]$$

- **Absorption by screened excitons** (Hulthen potential)
- Kramers-Kronig transform following Tanguy (includes broadening Γ)
- **Degenerate Fermi-Dirac statistics to calculate f_h and f_e .**
- Two terms for light and heavy excitons
- Non-parabolicity and temperature-dependent mass included from k.p theory
- **Screening parameter $g=12/\pi^2 a_R k_{TF}$ (large: no screening)**
- Only two free parameters: Band gap E_0 and broadening Γ
- Amplitude A and exciton binding energy R from k.p theory and effective masses

Conclusion

1. Optical and X-ray Characterization of thick **GeSn alloys** (1-2% Sn) **on GaAs**
Ken Hass Outstanding Student Paper Award March 2022, Forum on Industrial & Applied Physics, APS
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